

# The Jahn-Teller and pseudo-Jahn-Teller effects in propyne radical cation

Arun Kumar Kanakati, Vadala Jhansi Rani and S. Mahapatra

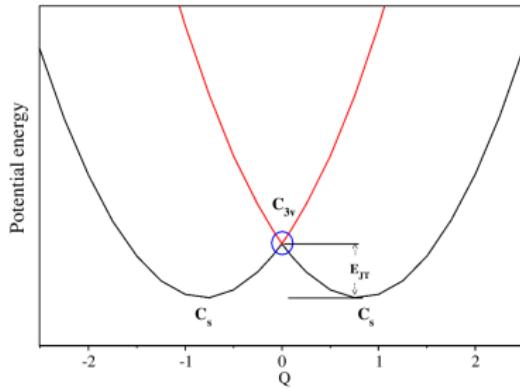
School of Chemistry  
University of Hyderabad  
India

May 15, 2023

# The Jahn-Teller theorem

- "... stability and degeneracy are not possible simultaneously unless the molecule is a linear one ..."<sup>1</sup>
- Structural aspects, global minimum of (JT or pseudo-JT) distorted potential energy surfaces<sup>2</sup>

- Break-down of the **Born-Oppenheimer** approximation
- Cls are crucial for signaling the ultrafast decay of excited molecular states
- Nonadiabatic situation



Schematic PES

<sup>1</sup>H. A. Jahn, E. Teller, and F. G. Donnan, Proc. Roy. Soc. A **161**, 220 (1937).

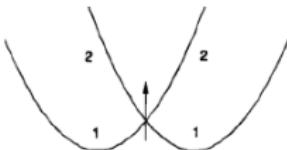
<sup>2</sup>I.B. Bersuker, The Jahn-Teller effect, Springer series in Chem. Phys. 2009.

# The Jahn-Teller theorem

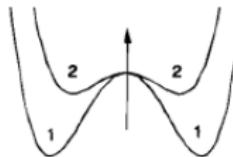
## Types of degeneracies

- ① Conical intersections : JT, PJT & accidental degeneracies
- ② Glancing interactions : RT

### Conical intersections



### Glancing interactions



- Group theory : Symmetry selection rule:  $[\Gamma_{el}^2] \supset \Gamma_{vib}$

- Three representative cases :

Linear molecules	$[\mathbf{E}_k^2] = \sigma + \delta$	no JTE
Tetragonal point groups	$[\mathbf{E}^2] = \mathbf{a}_1 + \mathbf{b}_1 + \mathbf{b}_2$	$\mathbf{E} \otimes \mathbf{b}$ JTE
Trigonal point groups	$[\mathbf{E}^2] = \mathbf{a}_1 + \mathbf{e}$	$\mathbf{E} \otimes \mathbf{e}$ JTE

## Motivation:

- **C<sub>3</sub>H<sub>4</sub>**: Allene (H<sub>2</sub>CCCH<sub>2</sub>, D<sub>2d</sub>), Propyne (H<sub>3</sub>CCCH, C<sub>3v</sub>), and Cyclopropene (C<sub>2v</sub>) are three stable isomers.
- Allene and Propyne are important intermediates in combustion and astrochemistry, and they are cumulene series with odd number of carbon atoms.
- Allene: E $\otimes$ b JTE and (E $\otimes$ b)+E PJTE
- Propyne: E $\otimes$ e JTE and (E $\otimes$ e)+A<sub>1</sub> PJTE

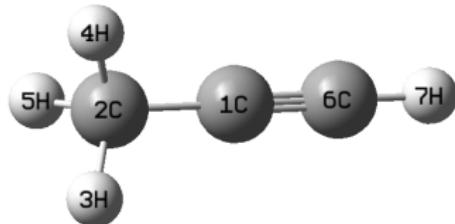
## Objectives of our work

- A detailed investigation of multi-mode JT and PJT effects in the first three,  $\tilde{X}(^2E)$ ,  $\tilde{A}(^2E)$ , and  $\tilde{B}(^2A_1)$  electronic states of  $\text{H}_3\text{CCCH}^{+}$
- Construct the vibronic model Hamiltonian using the standard vibronic coupling theory
- Extensive *ab initio* quantum chemistry calculations
- JT, PJT and spin-orbit coupling effects in the nuclear dynamics

### Electronic structure:

- Symmetry point group :  $C_{3v}$
- Opt/Freq : MP2/aug-cc-pVDZ, G09
- $\Gamma_{vib} = 5a_1 + 5e$
- The symmetrized direct product of  $E$  representations in the  $C_{3v}$  point group yields

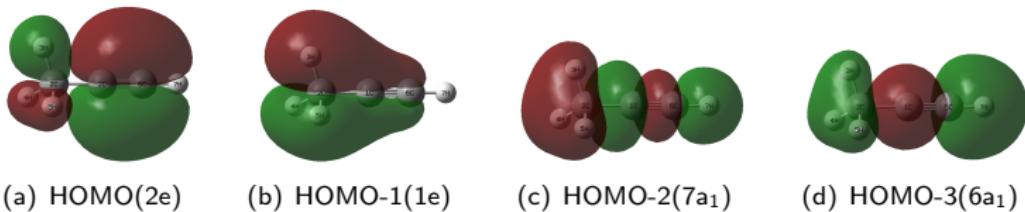
$$\begin{aligned}(E)^2 &= a_1 + e \\ E \otimes A_1 &= e\end{aligned}$$



Propyne equilibrium geometry

- The orbital configuration of propyne ( $\text{H}_3\text{CCCH}$ ) in its electronic ground state is

$$\tilde{X}(^1A_1) = (\text{core})(6a_1)^2(7a_1)^2(1e)^4(2e)^4$$



- Electronic configuration of first three low-lying electronic states of  $\text{H}_3\text{CCCH}^+$  is

$$\tilde{X}(^2E) = \dots(6a_1)^2(7a_1)^2(1e)^4(2e)^3$$

$$\tilde{A}(^2E) = \dots(6a_1)^2(7a_1)^2(1e)^3(2e)^4$$

$$\tilde{B}(^2A_1) = \dots(6a_1)^2(7a_1)^1(1e)^4(2e)^4$$

# Vertical ionization energies

States	$\tilde{X}(^2E)$	$\tilde{A}(^2E)$	$\tilde{B}(^2A_1)$	$\tilde{C}^2A_1$	
VIEs	10.23	15.03	15.30	17.78	OVGF
	10.40	15.09	15.31	17.68	EOMIP-CCSD
	10.28	15.40	15.75	18.29	CASSCF-MRC
	10.37	14.70	15.50/15.80	17.49	Expt. <sup>1</sup>
	10.37	14.4	15.13	17.2	Expt. <sup>2</sup>
	10.37	14.6	15.3/15.5	17.4	Expt. <sup>3</sup>
	10.37	13.69	15.2	17.2	Expt. <sup>4</sup>
	10.36	13.69/14.70/15.30	15.8	17.49	Expt. <sup>5</sup>

<sup>1</sup> M. H. Palmer, C. C. Ballard, and I. C. Walker, Chem. Phys. **249**, 129 (1999).

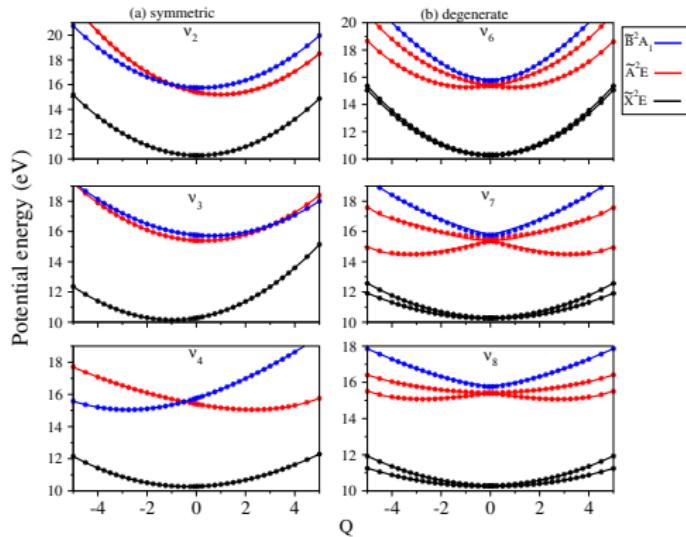
<sup>2</sup> C. Baker and D. W. Turner, Proc. R. Soc. London. Ser. A **308**, 19 (1968).

<sup>3</sup> W. Ensslin, H. Bock, and G. Becker, J. Am. Chem. Soc. **96**, 2757 (1974).

<sup>4</sup> D. Frost, F. Herring, C. McDowell, and I. Stenhouse, Chem. Phys. Lett. **4**, 533 (1970).

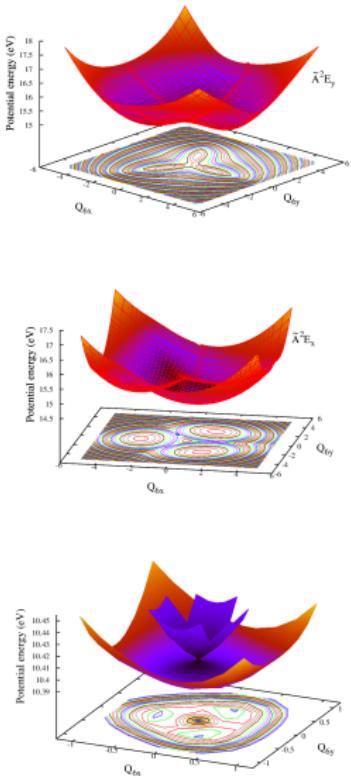
<sup>5</sup> G. H. Ho, M. S. Lin, Y. L. Wang, and T. W. Chang, J. Chem. Phys. **109**, 5868 (1998).

# Potential energy surfaces



- $E_{JT}(\tilde{X}(^2E)) = 0.02 \text{ eV} (\sim 160 \text{ cm}^{-1})$  and Expt.<sup>a</sup>,  $\sim 117 \text{ cm}^{-1}$
  - $E_{JT}(\tilde{A}(^2E)) = 1.14 \text{ eV} (\sim 11,372 \text{ cm}^{-1})$
  - Spin-orbit (SO) coupling of the  $\tilde{X}(^2E)$  state  $\sim -28.60 \text{ cm}^{-1}$  (EOMIP-CCSD) and  $\sim -28.83 \text{ cm}^{-1}$  (CASSCF-MRCI), and Marquez *et al.*<sup>a</sup> reported  $\sim -28 \text{ cm}^{-1}$

<sup>a</sup> D. R. Yarkony et al., J. Phys. Chem. A 117, 12002 (2013).



# Vibronic coupling model Hamiltonian

$$\mathcal{H} = (\mathcal{T}_N + \mathcal{V}_0)\mathbf{1}_5 + \Delta\mathcal{H}$$

$$\mathcal{T}_N = -\frac{1}{2} \sum_{i \in \mathfrak{a}_1} \omega_i \frac{\partial^2}{\partial Q_i^2} - \frac{1}{2} \sum_{i \in e} \omega_i \left( \frac{\partial^2}{\partial Q_{ix}^2} + \frac{\partial^2}{\partial Q_{iy}^2} \right),$$

$$\mathcal{V}_0 = \frac{1}{2} \sum_{i \in \mathfrak{a}_1} \omega_i Q_i^2 + \frac{1}{2} \sum_{i \in e} \omega_i (Q_{ix}^2 + Q_{iy}^2),$$

$$\Delta\mathcal{H} = \begin{pmatrix} u_{\tilde{X}}^x & u_{\tilde{X}}^{xy} & 0 & 0 & 0 \\ u_{\tilde{X}}^{xy} & u_{\tilde{X}}^y & 0 & 0 & 0 \\ 0 & 0 & u_A^x & u_A^{xy} & \sum_{i \in e} \lambda'_i Q_{ix} \\ h.c. & u_A^{xy} & u_A^y & \sum_{i \in e} -\lambda'_i Q_{iy} & u_{\tilde{B}}^y \end{pmatrix}.$$

diagonal:

$$\begin{aligned} u_j^{x/y} = & E_j^0 + \sum_{i \in \mathfrak{a}_1} \kappa_i^j Q_i + \frac{1}{2!} \sum_{i \in \mathfrak{a}_1} \gamma_i^j Q_i^2 + \frac{1}{3!} \sum_{i \in \mathfrak{a}_1} \sigma_i^j Q_i^3 + \frac{1}{4!} \sum_{i \in \mathfrak{a}_1} \delta_i^j Q_i^4 \\ & + \frac{1}{2!} \sum_{i \in e} \gamma_i^j (Q_{ix}^2 + Q_{iy}^2) + \frac{1}{3!} \sum_{i \in e} \sigma_i^j (2Q_{ix}^3 - 6Q_{ix}Q_{iy}^2) + \frac{1}{4!} \sum_{i \in e} \delta_i^j (Q_{ix}^4 + 2Q_{ix}^2Q_{iy}^2 + Q_{iy}^4) \\ & + \frac{1}{5!} \sum_{i \in e} \rho_i^j (2Q_{ix}^5 - 4Q_{ix}^3Q_{iy}^2 - 6Q_{ix}Q_{iy}^4) \pm \sum_{i \in e} \lambda_i^j Q_{ix} \pm \frac{1}{2!} \sum_{i \in e} \eta_i^j (Q_{ix}^2 - Q_{iy}^2) \\ & \pm \frac{1}{3!} \sum_{i \in e} \sigma_i'^j (Q_{ix}^3 + Q_{ix}Q_{iy}^2) \pm \frac{1}{4!} \sum_{i \in e} \delta_i'^j (Q_{ix}^4 - 6Q_{ix}^2Q_{iy}^2 + Q_{iy}^4) \pm \frac{1}{4!} \sum_{i \in e} \delta_i''^j (Q_{ix}^4 - Q_{iy}^4) \\ & \pm \frac{1}{5!} \sum_{i \in e} \rho_i'^j (Q_{ix}^5 - 10Q_{ix}^3Q_{iy}^2 + 5Q_{ix}Q_{iy}^4) \pm \frac{1}{5!} \sum_{i \in e} \rho_i''^j (Q_{ix}^5 + 2Q_{ix}^3Q_{iy}^2 + Q_{ix}Q_{iy}^4); j \in \tilde{X}, \tilde{A} \end{aligned}$$

# Vibronic coupling model Hamiltonian

$$u_j = E_j^0 + \sum_{i \in a_1} \kappa_i^j Q_i + \frac{1}{2!} \sum_{i \in a_1} \gamma_i^j Q_i^2 + \frac{1}{3!} \sum_{i \in a_1} \sigma_i^j Q_i^3 + \frac{1}{2!} \sum_{i \in e} \gamma_i^j (Q_{ix}^2 + Q_{iy}^2) + \frac{1}{4!} \sum_{i \in e} \delta_i^j (Q_{ix}^4 + 2Q_{ix}^2 Q_{iy}^2 + Q_{iy}^4)$$

off-diagonal:

$$\begin{aligned} u_j^{xy} = & \sum_{i \in e} \lambda_i^j Q_{iy} - \frac{1}{2!} \sum_{i \in e} 2\eta_i^j Q_{ix} Q_{iy} + \frac{1}{3!} \sum_{i \in e} \sigma_i'^j (Q_{ix}^2 Q_{iy} + Q_{iy}^3) \\ & + \frac{1}{4!} \sum_{i \in e} \delta_i'^j (4Q_{ix}^3 Q_{iy} - 4Q_{ix} Q_{iy}^3) + \frac{1}{4!} \sum_{i \in e} \delta_i''^j (-2Q_{ix}^3 Q_{iy} - 2Q_{ix} Q_{iy}^3) \\ & + \frac{1}{5!} \sum_{i \in e} \rho_i'^j (-5Q_{ix}^4 Q_{iy} + 10Q_{ix}^2 Q_{iy}^3 - Q_{iy}^5) + \frac{1}{5!} \sum_{i \in e} \rho_i''^j (Q_{ix}^4 Q_{iy} + 2Q_{ix}^2 Q_{iy}^3 + Q_{iy}^5); j \in \tilde{x}, \tilde{A} \end{aligned}$$

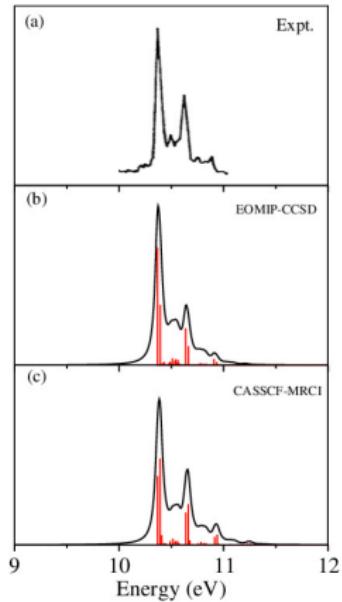
$$\begin{aligned} u_{AB}^x = & \sum_{i \in e} \lambda_i'^{(1)} Q_{ix} + \frac{1}{2!} \sum_{i \in e} \lambda_i'^{(2)} (Q_{ix}^2 - Q_{iy}^2) + \frac{1}{3!} \sum_{i \in e} \lambda_i'^{(3)} (Q_{ix}^3 + Q_{ix} Q_{iy}^2) \\ & + \frac{1}{4!} \sum_{i \in e} \lambda_i'^{(4)} (Q_{ix}^4 - 6Q_{ix}^2 Q_{iy}^2 + Q_{iy}^4) + \frac{1}{4!} \sum_{i \in e} \lambda_i'^{(4')} (Q_{ix}^4 - Q_{iy}^4) \\ & + \frac{1}{5!} \sum_{i \in e} \lambda_i'^{(5)} (Q_{ix}^5 - 10Q_{ix}^3 Q_{iy}^2 + 5Q_{ix} Q_{iy}^4) + \frac{1}{5!} \sum_{i \in e} \lambda_i'^{(5')} (Q_{ix}^5 + 2Q_{ix}^3 Q_{iy}^2 + Q_{ix} Q_{iy}^4), \end{aligned}$$

$$\begin{aligned} u_{AB}^y = & \sum_{i \in e} \lambda_i'^{(1)} Q_{iy} - \frac{1}{2!} \sum_{i \in e} 2\lambda_i'^{(2)} Q_{ix} Q_{iy} + \frac{1}{3!} \sum_{i \in e} \lambda_i'^{(3)} (Q_{ix}^2 Q_{iy} + Q_{iy}^3) \\ & + \frac{1}{4!} \sum_{i \in e} \lambda_i'^{(4)} (4Q_{ix}^3 Q_{iy} - 4Q_{ix} Q_{iy}^3) + \frac{1}{4!} \sum_{i \in e} \lambda_i'^{(4')} (-2Q_{ix}^3 Q_{iy} - 2Q_{ix} Q_{iy}^3) \\ & + \frac{1}{5!} \sum_{i \in e} \lambda_i'^{(5)} (-5Q_{ix}^4 Q_{iy} + 10Q_{ix}^2 Q_{iy}^3 - Q_{iy}^5) + \frac{1}{5!} \sum_{i \in e} \lambda_i'^{(5')} (Q_{ix}^4 Q_{iy} + 2Q_{ix}^2 Q_{iy}^3 + Q_{iy}^5). \end{aligned}$$

# Vibrational energy level spectrum:

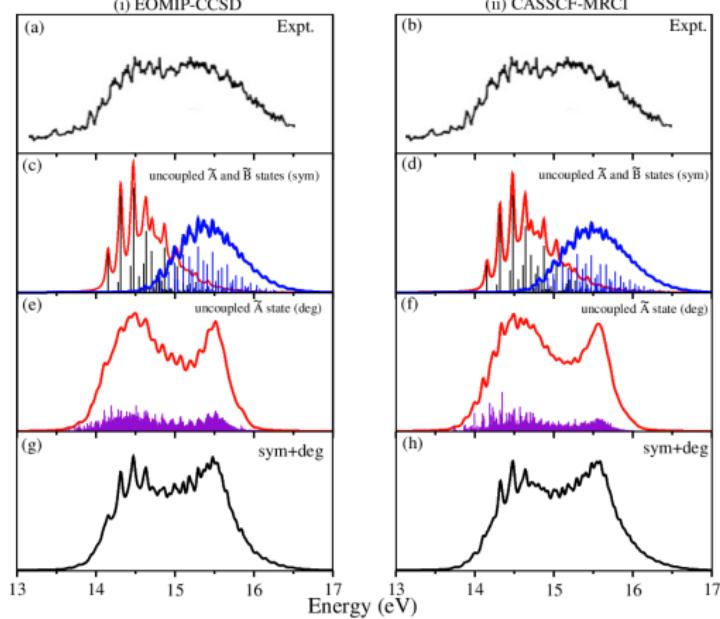
$\tilde{X}(^2E)$  state

Relative Intensity (arb. units)



$\tilde{A}(^2E)$ - $\tilde{B}(^2A_1)$  states

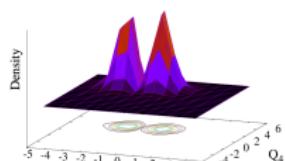
Relative intensity (arb. units)



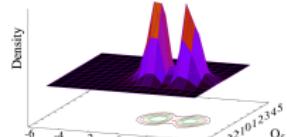
<sup>1</sup>C. Baker and D. W. Turner, Pro. R. Soc. London. Ser. A **308**, 19 (1968).

## Energy eigenvalues (EOMIP-CCSD)

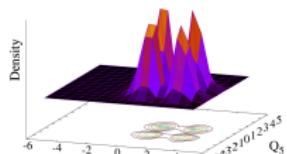
$\tilde{X}(^2E)$				$\tilde{A}(^2E)$			$\tilde{B}(^2A_1)$	
Energy	Ref. [1]	Ref. [2]	Assignment	Energy	Ref. [1]	Assignment	Energy	Assignment
0			$0^0$	0		$0^0$	0	$0^0$
876	940	$930 \pm 50$	$\nu_{50}^1$	1024		$\nu_{50}^1$	917	$\nu_{50}^1$
1328			$\nu_{40}^1$	1304	1290	$\nu_{40}^1$	1497	$\nu_{40}^1$
1752			$\nu_{50}^2$	2048		$\nu_{50}^2$	1835	$\nu_{50}^2$
2169	1940	$2000 \pm 50$	$\nu_{30}^1$	2145		$\nu_{30}^1$	2018	$\nu_{30}^1$
2204			$\nu_{40}^1 + \nu_{50}^1$	2328		$\nu_{40}^1 + \nu_{50}^1$	2414	$\nu_{40}^1 + \nu_{50}^1$
2628			$\nu_{50}^3$	2602		$\nu_{40}^2$	2752	$\nu_{50}^3$
2655			$\nu_{40}^2$	3073		$\nu_{50}^3$	2936	$\nu_{30}^1 + \nu_{50}^1$
3046			$\nu_{30}^1 + \nu_{50}^1$	3169		$\nu_{30}^1 + \nu_{50}^1$	2991	$\nu_{40}^2$
3067			$\nu_{20}^1$	3178		$\nu_{20}^1$	3017	$\nu_{20}^1$
3080			$\nu_{40}^1 + \nu_{50}^2$	3353		$\nu_{40}^1 + \nu_{50}^2$	3332	$\nu_{40}^1 + \nu_{50}^2$



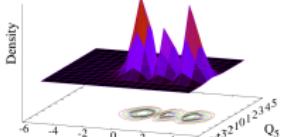
(e)  $1024 \text{ cm}^{-1}$



(f)  $1304 \text{ cm}^{-1}$



(g)  $2328 \text{ cm}^{-1}$



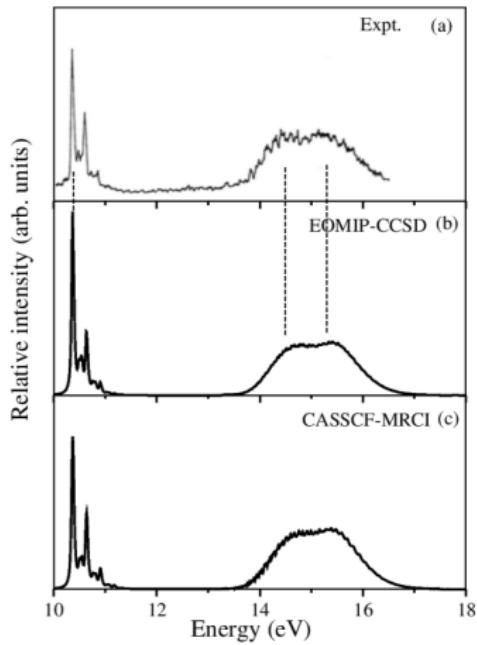
(h)  $2602 \text{ cm}^{-1}$

<sup>1</sup>C. Baker and D. W. Turner, Pro. R. Soc. London. Ser. A **308**, 19 (1968).

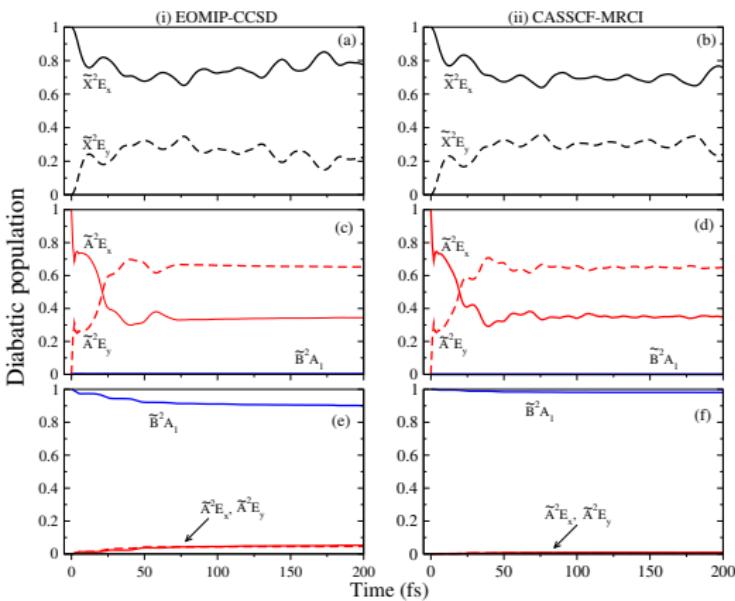
<sup>2</sup>U. Jacobella and F. Merkt, Mol. Phys. **116**, 302 (2018).

# Coupled state results

## Photoelectron spectrum of propyne



## Internal conversion dynamics



<sup>1</sup>C. Baker and D. W. Turner, Pro. R. Soc. London. Ser. A 308, 19 (1968).

## Conclusions:

- It is found that the  $\tilde{X}(^2E)$  electronic state is energetically well separated from the  $\tilde{A}(^2E)$  and  $\tilde{B}(^2A_1)$  states at the Franck-Condon geometry.
- Similar results were obtained for the allene radical cation.
- JT effect in the  $\tilde{A}(^2E)$  state is stronger than the  $\tilde{X}(^2E)$  state.
- The PJT coupling between the  $\tilde{A}(^2E)$ - $\tilde{B}(^2A_1)$  electronic states is weak in propyne. But it is stronger in the allene radical cation.
- In case of propyne, the complex band structure of the  $\tilde{A}(^2E)$  and  $\tilde{B}(^2A_1)$  states arises solely from their energetic proximity.
- This is in contrast to the effects in the vibronic band structure of  $\tilde{A}(^2E)$ - $\tilde{B}(^2B_2)$  electronic states of isomeric allene radical cation.

## Acknowledgement:

- Prof. Susanta Mahapatra
- Labmates, friends and family members
- UoH and CSIR
- 25th International conference on JT Effect organizers

# Thank You