



Beyond Born–Oppenheimer Constructed Diabatic Potential Energy Surfaces for HeH_2^+

Koushik Naskar¹, Satyam Ravi^{1,2}, Satrajit Adhikari^{1,*},
Michael Baer³, and Narayanasami Sathyamurthy⁴

¹School of Chemical Sciences, Indian Association for the Cultivation of Science
2A & 2B Raja S. C. Mullick Road, Jadavpur, Kolkata 700032, West Bengal, India

²School of Advanced Sciences and Languages VIT Bhopal University, Bhopal, 466114, India

³The Fritz Haber Center for Molecular Dynamics, The Hebrew University of Jerusalem, Jerusalem 91904, Israel

⁴Indian Institute of Science Education and Research Mohali, SAS Nagar, Manauli 140306, India

*E-mail: [pcsas@iacs.res.in](mailto:pcsa@iacs.res.in)

Poster presenter's e-mail: pckn@iacs.res.in

Beyond Born-Oppenheimer (BBO) Theory

Molecular Schrodinger Equation:

$$\hat{H}(e, n) |\Psi(e, n)\rangle = E |\Psi(e, n)\rangle$$

Molecular Hamiltonian:

$$\hat{H}(e, n) = \hat{T}_n(n) + \hat{H}_e(e; n)$$

Total molecular wave function:

$$\Psi(e, n) = \sum_{i=1}^N \psi_i(n) \xi_i(e; n)$$

Electronic Schrodinger Equation:

$$\hat{H}_e(e; n) \xi_i(e; n) = u_i(n) \xi_i(e; n)$$

Schrodinger Equation in Adiabatic Representation:

Non-Adiabatic Coupling Terms (NACTs)

NACTs from Hellmann-Feynman Theorem:

Cauchy's Residue Theorem

$$\int_0^{2\pi} \tau_{ij}^\phi(\phi) d\phi = n\pi$$

$$\left[-\frac{1}{2} \left(\vec{\nabla} + \vec{\tau} \right)^2 + (u - E) \right] \psi = 0$$

$$\vec{\tau}_{ij} = \left\langle \xi_i \left| \vec{\nabla} \right| \xi_j \right\rangle$$

$$\vec{\tau}_{ij} = \frac{\left\langle \xi_i \left| \vec{\nabla}_n \hat{H}_e \right| \xi_j \right\rangle}{(u_j - u_i)}$$

1. H. Hellmann, *Einfuhrang in die Quantenchemie* (Franz Duetche, Leipzig, 1937)
2. R. Feynman, *Phys. Rev.* **56**, 340 (1939); S. T. Epstein, *Am. J. Phys.* **22**, 613 (1954)
3. A.L. Cauchy, "Oeuvres complètes, Ser. 1", Paris **4** (1890)

Adiabatic to Diabatic Transformation (ADT):

Transforming the nuclear wave function

$$\psi(s_n) = A(s_n) \phi^d(s_n)$$

The ADT matrix, A satisfies the **ADT condition**

$$\vec{\nabla}_n A + \vec{\tau} A = \mathbf{0}$$

Schrodinger Equation in diabatic representation,

$$\left[-\frac{1}{2} \nabla^2 + (W - E) \right] \phi^d = \mathbf{0}$$

Diabatic Potential Energy (PE) matrix,

$$W = A^\dagger u A$$

□ The adiabatic-to-diabatic transformation matrix, A is orthogonal

➤ A matrix for a 2 state problem

$$\begin{pmatrix} \cos \Theta_{12} & \sin \Theta_{12} \\ -\sin \Theta_{12} & \cos \Theta_{12} \end{pmatrix} \quad \Theta_{12} \text{ is the ADT angle}$$

$$A(\Theta_{12}, \Theta_{13}, \Theta_{23}) = A_{12}(\Theta_{12}) \cdot A_{13}(\Theta_{13}) \cdot A_{23}(\Theta_{23})$$

➤ For three state problem, three Euler like ADT angles are required to construct the ADT matrix:

$$A = \begin{pmatrix} \cos \Theta_{12} & \sin \Theta_{12} & \mathbf{0} \\ -\sin \Theta_{12} & \cos \Theta_{12} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & 1 \end{pmatrix} \cdot \begin{pmatrix} \cos \Theta_{13} & \mathbf{0} & \sin \Theta_{13} \\ \mathbf{0} & 1 & \mathbf{0} \\ -\sin \Theta_{13} & \mathbf{0} & \cos \Theta_{13} \end{pmatrix} \cdot \begin{pmatrix} 1 & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \cos \Theta_{23} & \sin \Theta_{23} \\ \mathbf{0} & -\sin \Theta_{23} & \cos \Theta_{23} \end{pmatrix}$$

□ Features of Adiabatic to Diabatic Transformation:

- For a N dimensional electronic manifold, the ADT matrix can be considered as a product of elementary rotation matrices constituted with mixing angles between any two electronic states for the N -state sub-Hilbert space (SHS).

$$A = P_n \{ A^{12}(\Theta_{12}).A^{13}(\Theta_{13}).A^{23}(\Theta_{23}).....A^{N-1,N}(\Theta_{N-1,N}) \}, \quad n=1,...\Lambda!; \quad \Lambda = {}^N C_2 = \frac{N(N-1)}{2}$$

- Elementary rotation matrices A^{mn} are given by :

$$[A^{mn}(\Theta_{mn})]_{ij} = \delta_{ij}; \quad \{i, j\} \neq \{m, n\}.$$

$$A^{14}(\Theta_{14}) = \begin{pmatrix} \cos \Theta_{14} & 0 & 0 & \sin \Theta_{14} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\sin \Theta_{14} & 0 & 0 & \cos \Theta_{14} \end{pmatrix}$$

$$[A^{mn}(\Theta_{mn})]_{mm} = \cos \Theta_{mn} = [A^{mn}(\Theta_{mn})]_{nn}; \quad m \neq n$$

$$[A^{mn}(\Theta_{mn})]_{mn} = \sin \Theta_{mn} = -[A^{mn}(\Theta_{mn})]_{nm}; \quad m \neq n$$

- Using the ADT condition we get Λ unique coupled differential equation in the form:

$$\vec{\nabla}_n A + \vec{\tau} A = 0$$



$$\vec{\nabla}_n \Theta_{ij} = \sum_{m=1}^{\Lambda} c^{(m)} \vec{\tau}_{(m)}$$

1. M. Baer, *Chem. Phys. Letters*, **35**, 112 (1975)
2. M. Baer, *Chem. Phys. Letters*, **35**, 112 (1975).
3. Z. H. Top and M. Baer, *J. Chem. Phys.* **66**, 1363 (1977).
4. K. Naskar, S. Mukherjee, B. Mukherjee, S. Ravi, S. Mukherjee, S. Sardar and S. Adhikari, *J. Chem. Theory Comput.*, **16**, 1666–1680 (2020).

Application of BBO Theory for Scattering Processes: Ab - initio PESs, NACTs and Diabatic PESs of HeH₂⁺ system

- ❖ The *ab-initio* calculations on the four lowest singlet states of HeH₂⁺ ($1^1A'$, $2^1A'$, $3^1A'$ and $4^1A'$) are performed by **MOLPRO - 2018** quantum chemistry package in hyperspherical coordinate as a function hyperangles (θ, φ) for fixed values of hyperradius (ρ)
- ❖ MCSCF calculation is performed with three (3) electrons distributed over eight (8) orbitals followed by MRCI calculation using cc- pVQZ basis set.

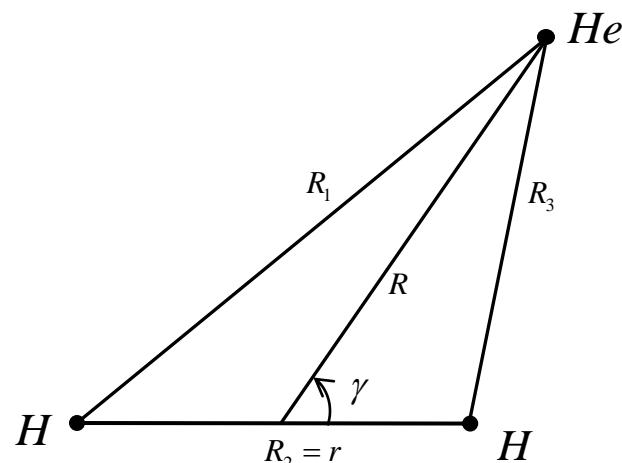
$$R_1 = \frac{\rho}{\sqrt{2}} d_3 [1 + \sin \theta \cos(\varphi + \varepsilon_3)]^{1/2},$$

$$R_2 = \frac{\rho}{\sqrt{2}} d_1 [1 + \sin \theta \cos \varphi]^{1/2},$$

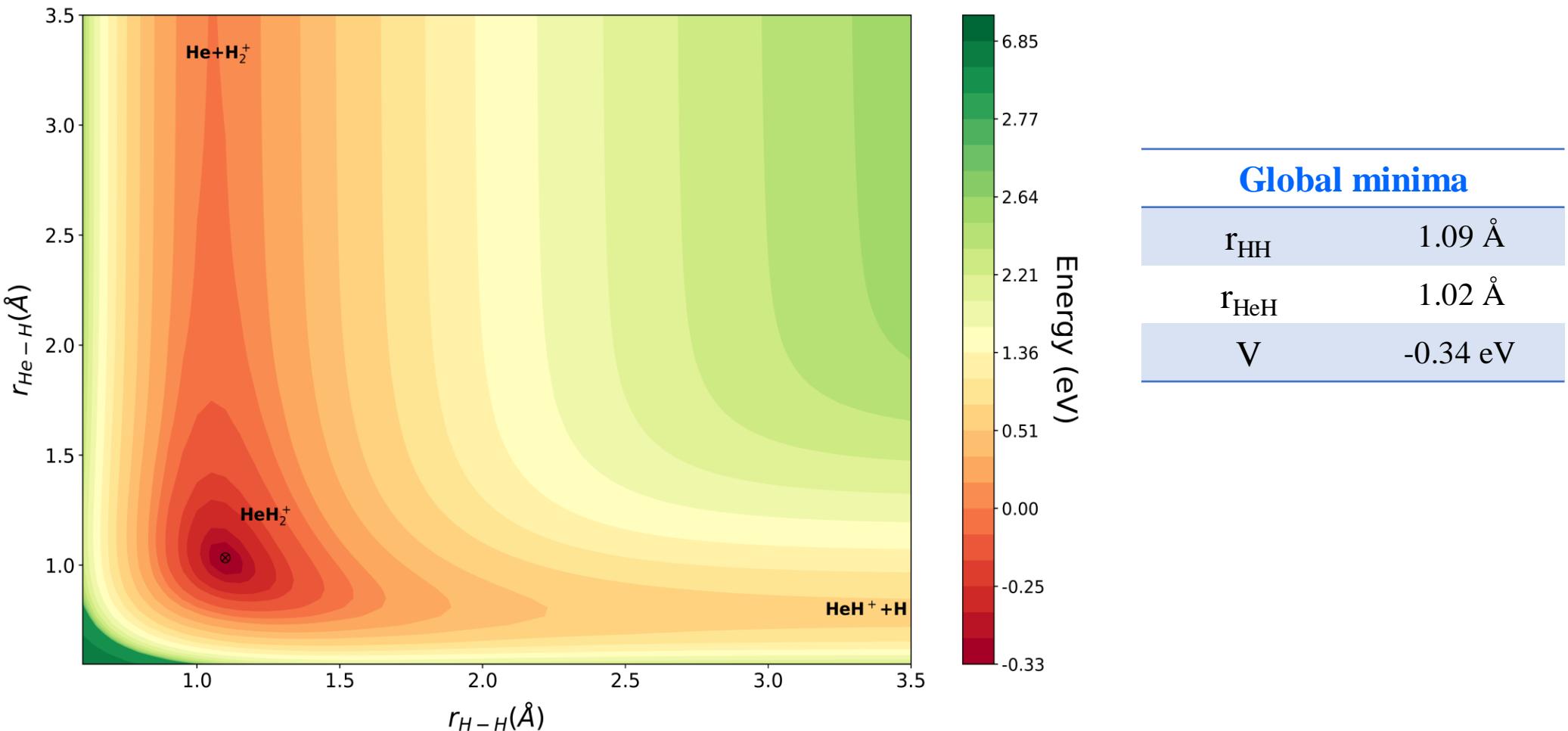
$$R_3 = \frac{\rho}{\sqrt{2}} d_2 [1 + \sin \theta \cos(\varphi - \varepsilon_2)]^{1/2}.$$

$$d_i = \sqrt{m_i(m_j + m_k) / \mu M} \quad \mu = \sqrt{m_1 m_2 m_3 / M}$$

$$\varepsilon_3 = 2 \tan^{-1}(m_2 / \mu) \quad \varepsilon_3 = 2 \tan^{-1}(m_2 / \mu)$$



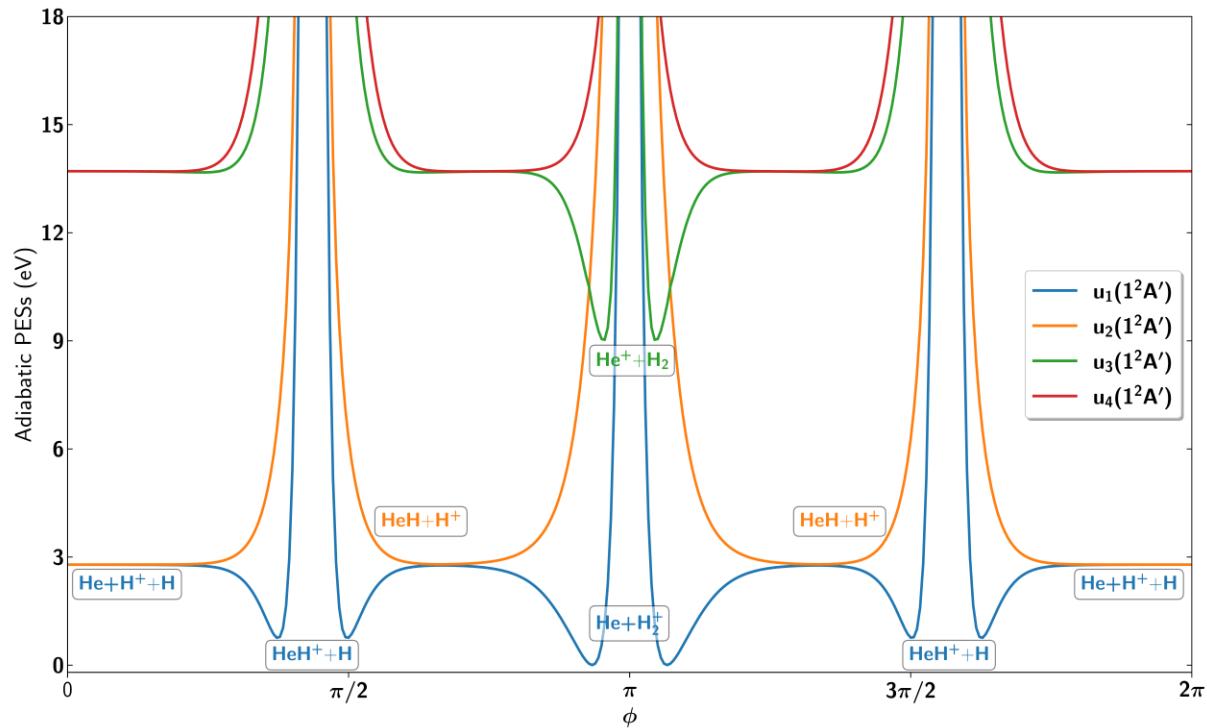
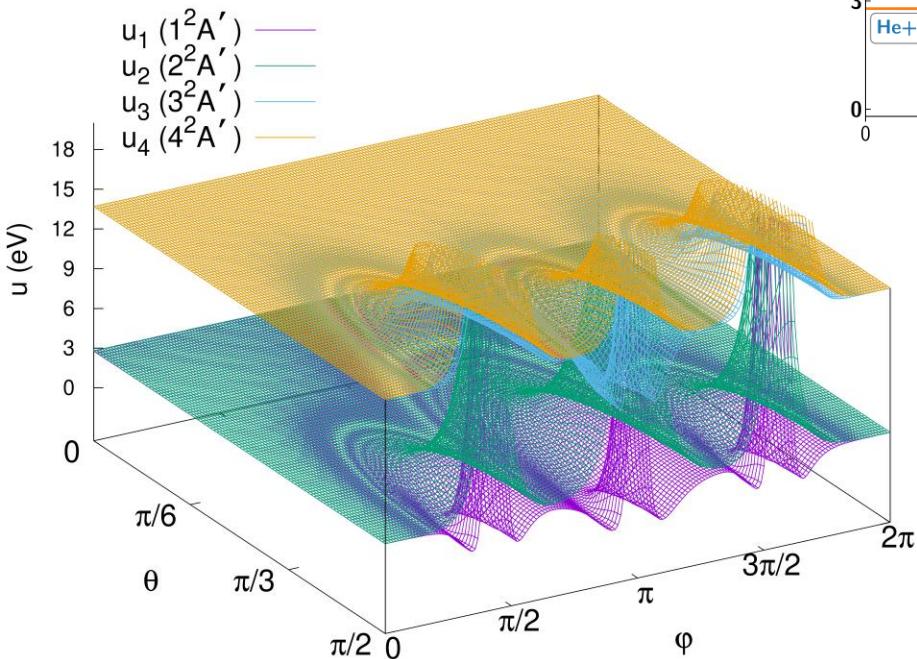
Adiabatic PESs of HeH_2^+ system



- Adiabatic ground electronic state of the HeH_2^+ system in collinear geometry

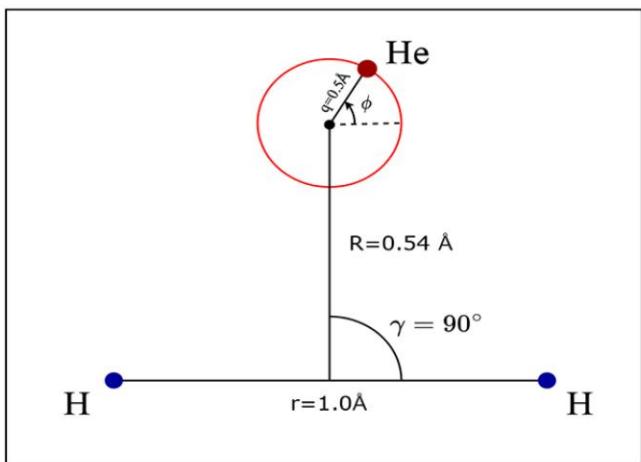
Adiabatic PESs of HeH_2^+ system

Different channels for the HeH_2^+ system at asymptotic ρ

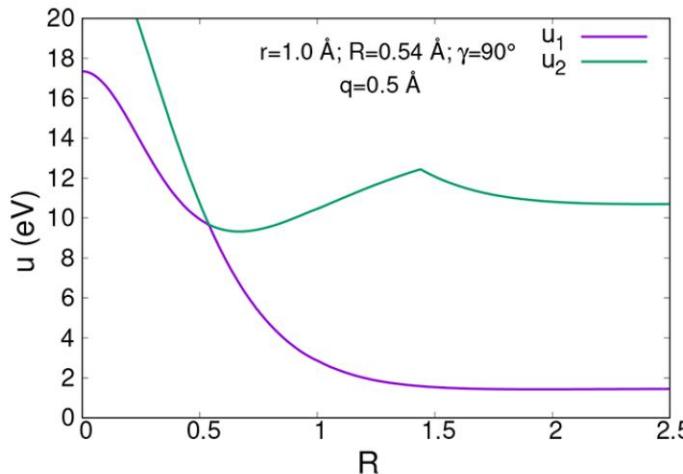


Adiabatic PESs for the HeH_2^+ system
in hyperspherical coordinate

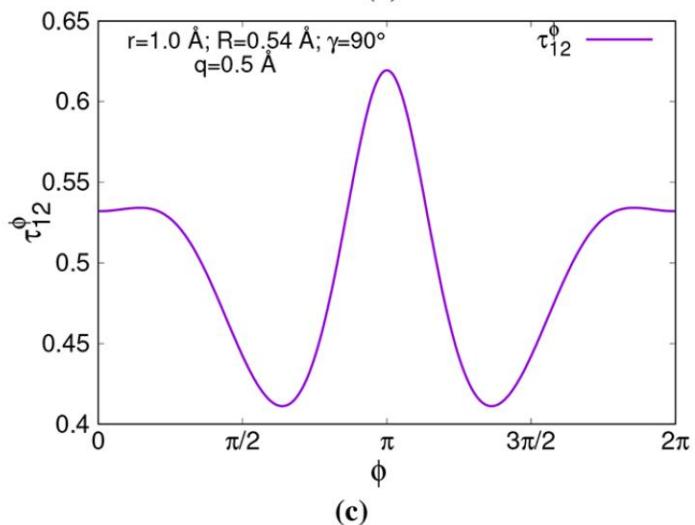
BBO Theory: Locating the conical intersections



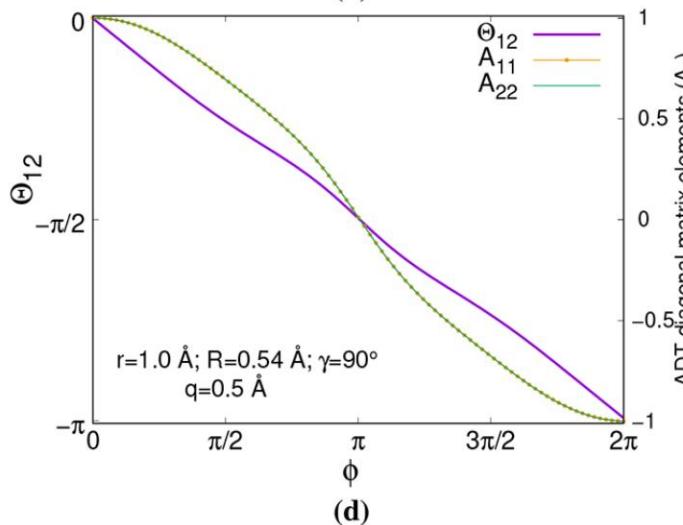
(a)



(b)



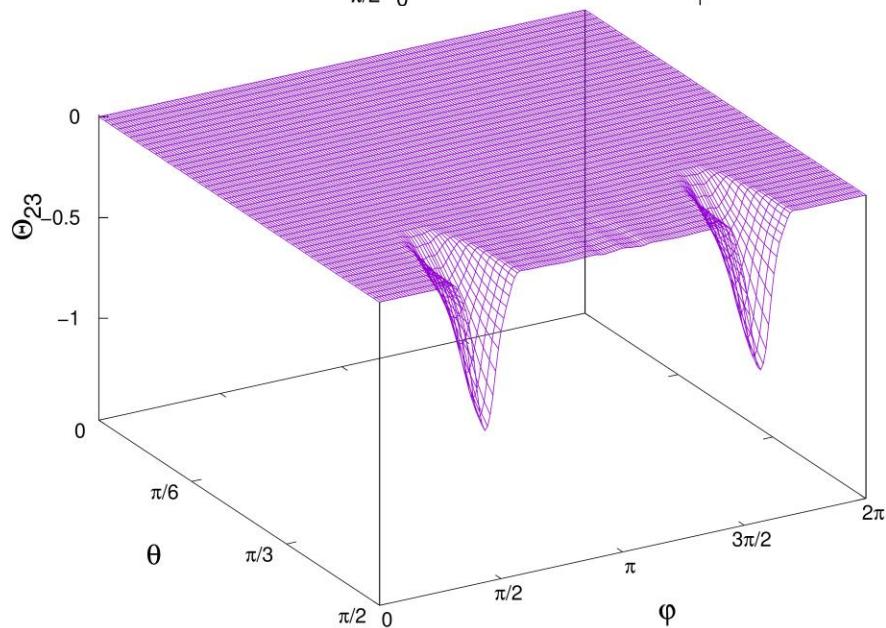
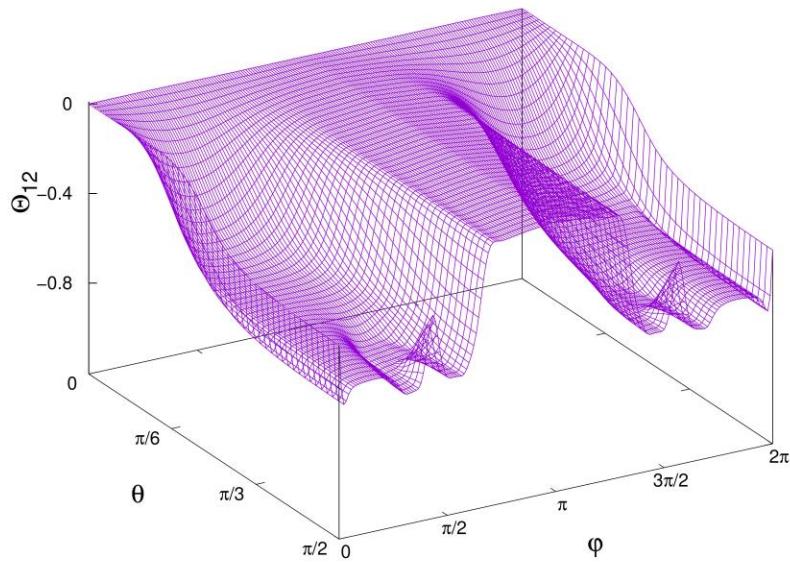
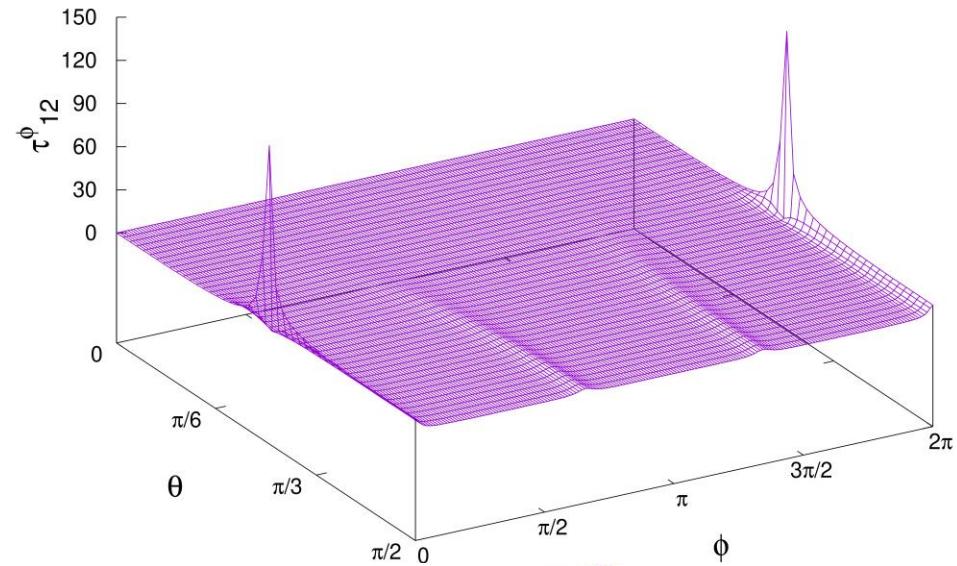
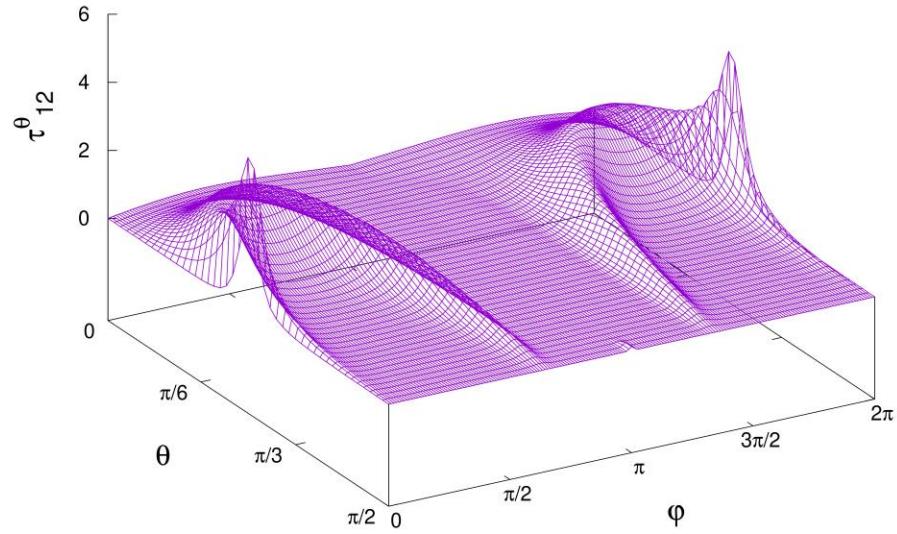
(c)



(d)

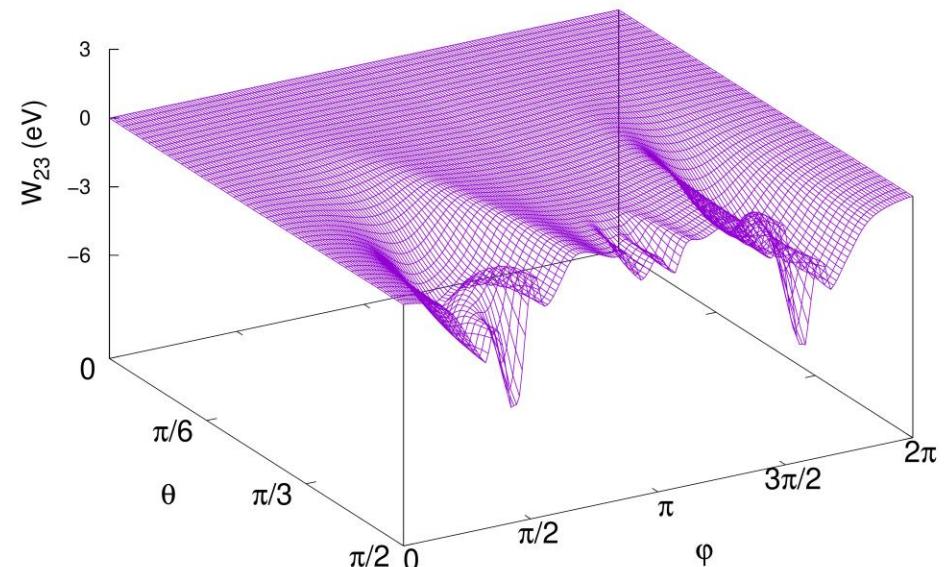
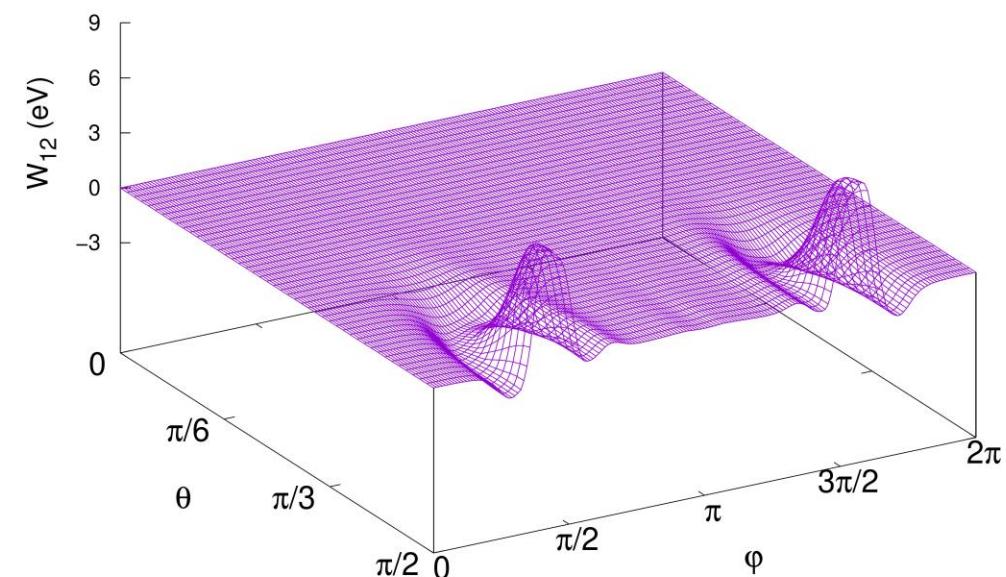
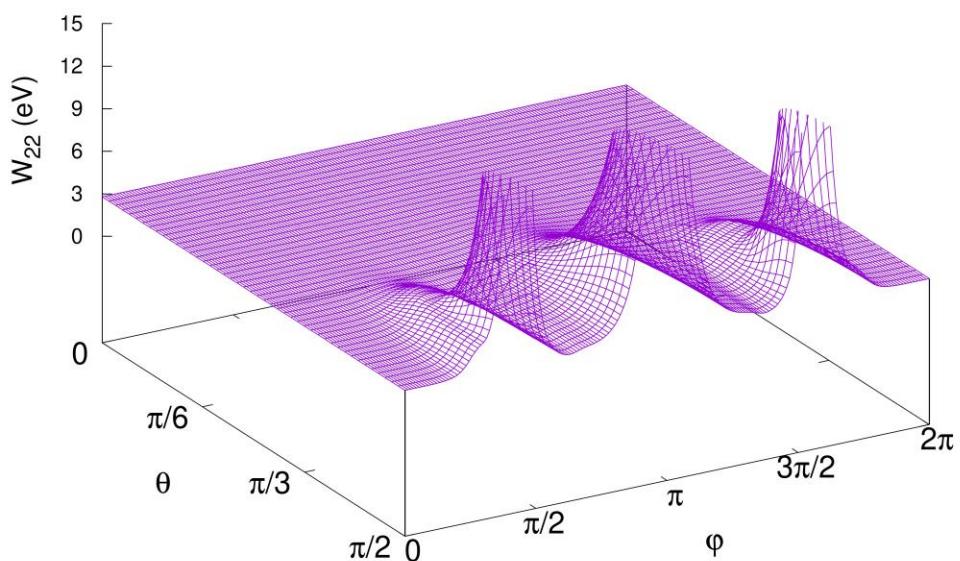
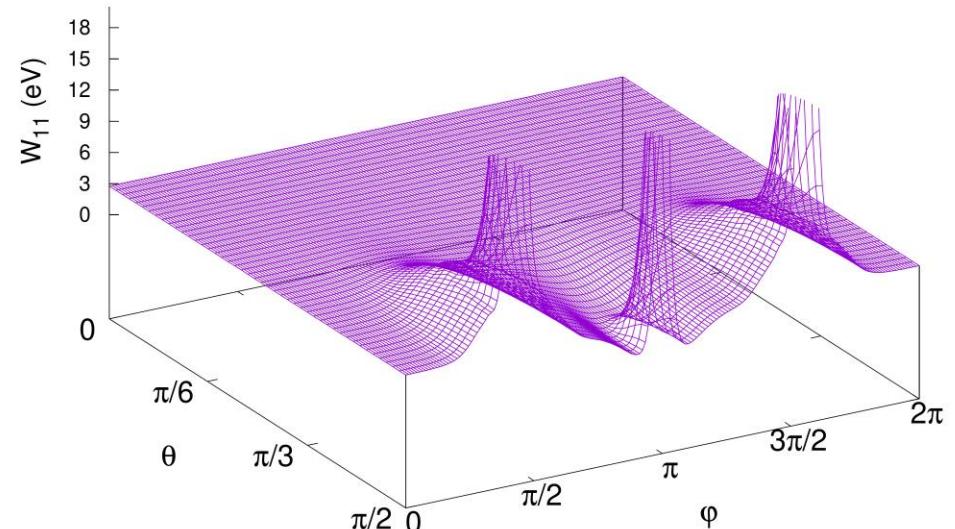
CI between ground and first excited state in C_{2v} configuration

NACTs and ADT Angles of HeH_2^+ system



1. K. Naskar, S. Ravi, S. Adhikari, M. Baer and N. Sathyamurthy; *J. Phys. Chem. A*, **127**, 3832 (2023)

Diabatic PESs and couplings of HeH_2^+ system



Summary & Future Plans

- Ab-initio based adiabatic PESs and NACTs are calculated for lowest four states of HeH_2^+ system.
- Conical intersection (CI) is located between the states in collinear as well as C_{2v} configurations is located and validated by integrating the NACTs along appropriately chosen contours where the resulting ADT angle reaches multiple integers of π .
- ADT angles are determined by solving the ADT equations to construct the diabatic potential matrix for the HeH_2^+ system which are smooth, single-valued, continuous, and symmetric.

Future Plans

- Utilize the four state diabatic potential matrix to study the ground state proton transfer reaction $\text{He} + \text{H}_2^+ \rightarrow \text{HeH}^+ + \text{H}$ as well as the excited state hydrogen transfer reaction $\text{He}^+ + \text{H}_2 \rightarrow \text{HeH}^+ + \text{H}$

Acknowledgment

K.N. thanks CSIR, India [Grant No. 09/080(1068)/2018-EMR-I] for funding. SR thanks VIT Bhopal for the seed grant, Support for Excellence in Academic Research (SPEAR), (File No: VITB/SEEDGRANT/2022/08). S.A. acknowledges the SERB (Project No. CRG/2019/000793) for the research funding. S.A. also acknowledges IACS for the CRAY super-computing facility.